# Band offset determination for amorphous Al<sub>2</sub>O<sub>3</sub> deposited on bulk AlN and atomic-layer epitaxial AlN on sapphire

Cite as: Appl. Phys. Lett. **117**, 182103 (2020); https://doi.org/10.1063/5.0025835 Submitted: 18 August 2020 . Accepted: 26 October 2020 . Published Online: 05 November 2020

D Chaker Fares, Fan Ren, Marko J. Tadjer, D Jeffrey Woodward, Michael A. Mastro, Boris N. Feigelson, Charles R. Eddy, and S. J. Pearton

### COLLECTIONS

Paper published as part of the special topic on Ultrawide Bandgap Semiconductors







### **ARTICLES YOU MAY BE INTERESTED IN**

Graphene-induced crystal-healing of AIN film by thermal annealing for deep ultraviolet lightemitting diodes

Applied Physics Letters 117, 181103 (2020); https://doi.org/10.1063/5.0028094

Full-composition-graded  $In_XGa_{1-X}N$  films grown by molecular beam epitaxy Applied Physics Letters 117, 182101 (2020); https://doi.org/10.1063/5.0021811

Electron and hole mobility of rutile GeO<sub>2</sub> from first principles: An ultrawide-bandgap semiconductor for power electronics

Applied Physics Letters 117, 182104 (2020); https://doi.org/10.1063/5.0033284





## Band offset determination for amorphous $Al_2O_3$ deposited on bulk AIN and atomic-layer epitaxial AIN on sapphire

Cite as: Appl. Phys. Lett. **117**, 182103 (2020); doi: 10.1063/5.0025835 Submitted: 18 August 2020 · Accepted: 26 October 2020 · Published Online: 5 November 2020







### **AFFILIATIONS**

- Department of Chemical Engineering, University of Florida, Gainesville, Florida 32611, USA
- <sup>2</sup>Naval Research Laboratory, Washington, DC 20375, USA
- <sup>3</sup>American Society for Engineering Education Postdoctoral Fellow, Residing at Naval Research Laboratory, Washington, DC 20375, USA

Note: This paper is part of the Special Topic on Ultrawide Bandgap Semiconductors.

a) Author to whom correspondence should be addressed: marko.tadjer@nrl.navy.mil

### **ABSTRACT**

Valence and conduction band offsets of atomic layer deposition (ALD)  $Al_2O_3$  deposited on bulk AlN crystals were determined using x-ray photoelectron spectroscopy to be  $\Delta E_V = 0.75$  eV and  $\Delta E_C = -1.45$  eV, with a measured energy gap of the  $Al_2O_3$  film of 6.9 eV. In addition, crystalline AlN deposited by atomic layer epitaxy on sapphire was evaluated, resulting in a valence band offset of  $\Delta E_V = -0.75$  eV and a conduction band offset of  $\Delta E_C = 3.25$  eV due to the wider bandgap of the crystalline  $Al_2O_3$  substrate compared to amorphous ALD  $Al_2O_3$ . Both heterojunctions exhibited type-II behavior and similar valence band offsets.

Published under license by AIP Publishing. https://doi.org/10.1063/5.0025835

The past several years have witnessed the emergence of a number of wide and ultra-wide bandgap (UWBG) semiconductors for electronic device applications. Gallium-based compounds such as gallium nitride (GaN), which has found widespread commercial application, and gallium oxide (Ga2O3), which has only recently gathered widespread interest as a device material, both offer advantageous properties such as a high critical field.<sup>2,3</sup> Additionally, the capability to incorporate Al and form ternary alloys has further extended the capabilities of these material systems into heteroepitaxial device structures such as AlGaN/GaN and (AlGa)<sub>2</sub>O<sub>3</sub>/Ga<sub>2</sub>O<sub>3</sub> heterostructure field effect transistors (HFETs).4,5 At high Al concentrations, AlN and Al2O3 are the widest bandgap III-Nitride and III-Oxide components of these material systems.<sup>6-9</sup> In particular, AlN technology has demonstrated native substrates, high quality epitaxial growth, and possibility for relatively shallow n-type donors. While AlN/GaN HFETs have been developed for some time, the first AlN channel transistor with implanted source/ drain was only recently demonstrated. 10-12 Most recently, improved understanding of Si doping in AlN has renewed interest in this UWBG material for electronic devices for high breakdown voltage,

high power applications. <sup>13–17</sup> As a result, a number of fundamental electronic experiments, such as the presently reported band diagrams, are needed for UWBG materials such as AlN, as has been the case with recent band diagram investigations for other promising materials such as  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. <sup>2</sup>

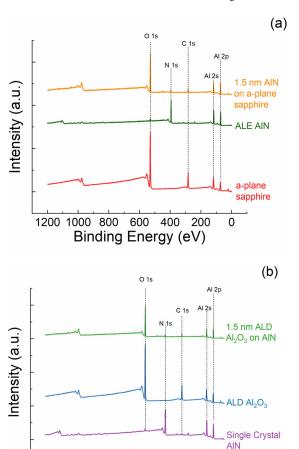
One of the fundamental methods of understanding the electronic properties of heterostructures between different materials is the construction of their band diagram. In the literature, band offset reports of oxide-semiconductor and semiconductor-semiconductor heterostructures abound as device technology has advanced over the years. Particularly for the III-N and III-O systems, a large body of literature exists for band offsets as different gate dielectrics are explored: SiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, and HfO<sub>2</sub> being among the most popular. Ternary nitride alloys with Al and their band structures have been reported. Band offsets between the binary III-O and III-N semiconductor heterojunctions have also been recently explored. The band alignments of AlN, GaN, and InN heterojunctions to Ga<sub>2</sub>O<sub>3</sub>, as well as GaN-Al<sub>2</sub>O<sub>3</sub> and InN-Al<sub>2</sub>O<sub>3</sub>, are all known. 19-23 The InN-In<sub>2</sub>O<sub>3</sub> heterojunction, which is relevant for thin-film transistor technology, has been studied as

Department of Materials Science and Engineering, University of Florida, Gainesville, Florida 32611, USA

well. <sup>24</sup> This paper aims to bridge the current gap in band offset knowledge between oxides and nitrides by characterizing the AlN-Al $_2$ O $_3$  heterojunction using deposited films of Al $_2$ O $_3$  and AlN on single-crystalline substrates of AlN and Al $_2$ O $_3$ , respectively.

We deposited amorphous Al<sub>2</sub>O<sub>3</sub> on bulk single-crystalline AlN substrates using atomic layer deposition. In addition, atomic-layer deposited epitaxial AlN has been grown on substrates of a-plane Al<sub>2</sub>O<sub>3</sub>. Al<sub>2</sub>O<sub>3</sub> was deposited by atomic layer deposition (ALD) on single-crystalline AlN substrates obtained commercially. These samples are, henceforth, referred to as Al<sub>2</sub>O<sub>3</sub>/AlN. Both thick (200 nm) and thin (1.5 nm) layers of the Al<sub>2</sub>O<sub>3</sub> dielectric were deposited to measure both the bandgaps and core levels. For pre-deposition substrate cleaning, the following rinse sequence was employed: acetone, isopropyl alcohol (IPA), dry N<sub>2</sub>, and finally ozone exposure for 15 min. ALD Al<sub>2</sub>O<sub>3</sub> was deposited at 200 °C in a Cambridge Nano Fiji 200 using a trimethylaluminum source and a remote inductively coupled plasma (ICP) of O<sub>2</sub> at 300 W.<sup>25,26</sup>

The atomic layer epitaxial (ALE) AlN films on the a-place sapphire, henceforth referred to as AlN/Al<sub>2</sub>O<sub>3</sub>, were grown at 525 °C



**FIG. 1.** XPS survey scans of (a) single crystal AlN, ALD  $Al_2O_3$ , and a heterostructure of  $Al_2O_3$  on AlN and (b) a-plane sapphire, ALE AlN, and a heterostructure of AlN on sapphire. The intensity is in arbitrary units (a.u.).

400

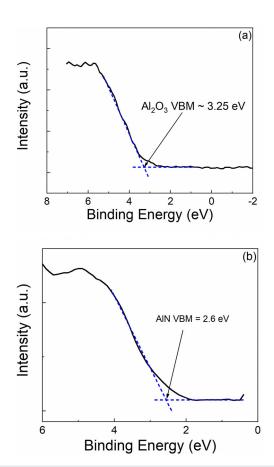
0

600

Binding Energy (eV)

using a Veeco/Cambridge Nanotech (CNT) Fiji 200 and ultra-high purity (UHP) gases, with constant 100 and 30 sccm flows of Ar through the plasma source and precursor line, respectively. Prior to AlN deposition, the  $Al_2O_3$  substrates were cleaned by sonicating in sequential 5 min baths of acetone and isopropanol at  $40\,^{\circ}$ C and a final bath in de-ionized water. The substrates were then dried in  $N_2$ , loaded into the reactor, and heated to growth temperature. An *in situ* cleaning process was performed immediately prior to growth. The *in situ* cleaning process consisted of subjecting the substrates to 10 cycles each of  $H_2$  and  $N_2$  plasma exposure at 300 W forward power, during which 75 (50) sccm  $H_2$  ( $N_2$ ) was flowed through the plasma source in addition to the Ar. Every cycle was 20 s in duration and was followed by a 10 s purge with only Ar flowing.

The AlN growth process consisted of 500 cycles of the following. First, the sample surface was saturated by a 60 ms pulse of trimethylaluminum. The chamber was then purged by the Ar flow for 8 s in order to remove excess precursor and reaction byproducts. Next, the precursor-terminated surface was exposed to a mixed  $N_2/H_2$  plasma at 300 W forward power for 20 s, during which 75 sccm of  $N_2$  and 10 sccm of  $H_2$  were flowed through the plasma source in addition to the Ar. Finally, residual gases and reaction byproducts were removed by a second Ar purge, completing the cycle.



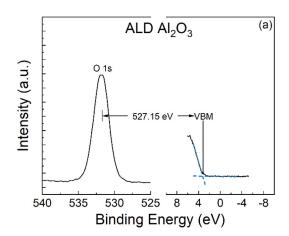
**FIG. 2.** XPS spectra of the valence band maxima (VBM) for (a) reference thick-film ALD  $Al_2O_3$  and (b) thick-film ALE AIN. The intensity is in arbitrary units (a.u.).

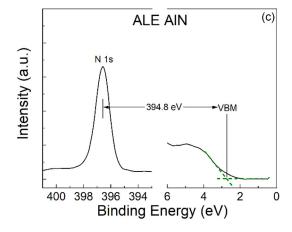
1200 1000 800

After deposition, the  $Al_2O_3/AlN$  and  $AlN/Al_2O_3$  samples were transferred into the chamber of an ULVAC PHI XPS system. A monochromatic, Al x-ray source with a power of 300 W and a dual beam charge neutralization system with simultaneous low-energy electron and ion means were used to perform the charge-compensated XPS measurement. Charge correction was performed using the adventitious carbon (C–C) line in the C 1s spectra at 284.8 eV. All spectroscopic equipment and electron analyzers were grounded. Samples were electronically insulated from the chuck to avoid differential charging effects, which were not observed. The bandgap of  $Al_2O_3$  was measured using reflection electron energy loss spectroscopy (REELS) with a 1 kV electron beam and the hemispherical electron analyzer.

The band offsets of the  $Al_2O_3/AlN$  and  $AlN/Al_2O_3$  heterojunctions were measured with high precision using the XPS method developed by Kraut *et al.*<sup>29</sup> To obtain band offsets using this technique, the core energy levels and the valence band maxima of each material need to be measured first. For the  $Al_2O_3/AlN$  heterojunction, the valence band offset (VBO) is then given by

$$\Delta E_V = \left(E_{core}^{Al2O3} - E_{VBM}^{Al2O3}\right) - \left(E_{core}^{AlN} - E_{VBM}^{AlN}\right) - \left(E_{core}^{O1s} - E_{core}^{N1s}\right). \tag{1}$$



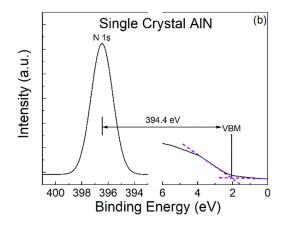


For the conduction band offset of the heterojunction, the bandgap of each constituent material needs to be known. The conduction band offset is thus given by

$$\Delta E_C = E_G^{Al2O3} - E_G^{AlN} - |\Delta E_V|. \tag{2}$$

Figure 1 shows the XPS scans of the two  $Al_2O_3$  films and the bulk AlN reference. The signature Al 2p, N 1s, and O 1s peaks from the two materials expected to be present were readily observable. An O 1s peak was observable on the AlN substrate reference scan [Fig. 1(a)] but was greatly suppressed and was likely due to O impurities in the crystal. The N 1s peak showed the highest intensity in the AlN substrate and was suppressed as thin ALD  $Al_2O_3$  was deposited. The N 1s peak was not observed from the sample with thick ALD  $Al_2O_3$  on the AlN substrate. Nearly identical peaks were observable from the XPS scans of the ALE AlN films on a-plane sapphire, as shown in Fig. 1(b).

Figure 2 shows the measured valence band maxima (VBM) from the thick  $Al_2O_3$  film [Fig. 2(a)] and the bulk AlN substrate [Fig. 2(b)] as 3.25 and 2.6 eV, respectively. Identical results (not shown) were obtained for the AlN/sapphire sample. This is particularly interesting



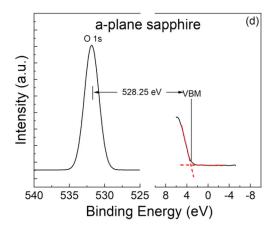


FIG. 3. High resolution XPS spectra for the vacuum-core delta regions of (a) reference thick film ALD  $Al_2O_3$  and (b) reference single crystal AIN for the ALD  $Al_2O_3$ /AIN sample; (c) thick film ALE AIN and (d) reference  $Al_2O_3$  for the ALE AIN/ $Al_2O_3$  sample. The intensity is in arbitrary units (a.u.).

TABLE I. Summary of measured core levels (eV) for single crystal AIN, ALD Al<sub>2</sub>O<sub>3</sub>, and a heterostructure of ALD Al<sub>2</sub>O<sub>3</sub> deposited on bulk AIN.

Single crystal AlN					A	LD Al <sub>2</sub> O <sub>3</sub>	Thin ALD Al <sub>2</sub> O <sub>3</sub> on single crystal AlN		
Core level	VBM	Core level peak	Core-VBM	Core level	VBM	Core level peak	Core—VBM	Δ Core level N1s—O1s	Valence band offset (VBO)
N1s	2.1	396.5	394.4	O1s	3.25	530.4	527.15	-133.5	$0.75 \pm 0.05$

TABLE II. Summary of measured core levels (eV) for a-plane sapphire, ALE AIN, and a heterostructure of AIN on sapphire.

Reference Al <sub>2</sub> O <sub>3</sub>					Refe	erence AlN	Thin AlN on Al <sub>2</sub> O <sub>3</sub>		
Core level	VBM	Core level peak	Core—VBM	Core level	VBM	Core level peak	Core—VBM	Δ Core level O1s—N1s	Valence band offset (VBO)
O1s	3.25	531.50	528.25	N1s	2.6	397.4	394.8	134.2	$-0.75 \pm 0.05$

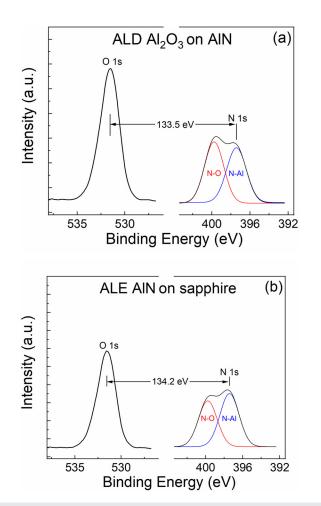
considering that amorphous ALD Al<sub>2</sub>O<sub>3</sub> and a bulk sapphire crystal were compared in this work. In comparison to previous work, similar valence band maxima for amorphous and crystalline Al<sub>2</sub>O<sub>3</sub> were also reported by Filatova and Konashuk<sup>31</sup> and Balzarotti and Bianconi.<sup>32</sup>

To obtain the energy difference between the VBM energies for each material and the core shell N 1s electrons for AlN and O 1s electrons for Al<sub>2</sub>O<sub>3</sub>, the previously measured VBM maxima are compared with core shell XPS data obtained from the 200 nm thick films. The energy difference spectra are shown in Fig. 3(a) for the ALD Al<sub>2</sub>O<sub>3</sub>, Fig. 3(b) for the bulk AlN substrate, Fig. 3(c) for the ALE AlN film, and Fig. 3(d) for the sapphire substrate. The energies measured from each comparison were 527.15 eV [Fig. 3(a)], 394.4 eV [Fig. 3(b)], 394.8 eV [Fig. 3(c)], and 528.25 eV [Fig. 3(d)], respectively. These results are summarized in the first two columns of Tables I and II.

To complete obtaining the experimental data required in Eq. (1), the energy difference between the same two core shell electron energies at the interface of the heterojunction was measured on samples of a 1.5 nm thick film on a bulk substrate. Figure 4(a) shows the core level difference from the 1.5 nm thick ALD Al<sub>2</sub>O<sub>3</sub> film on the AlN substrate to be 133.5 eV. For the 1.5 nm ALE AlN film on sapphire, this energy difference was measured to be 134.2 eV, as shown in Fig. 4(b). Thus, according to Eq. (1), a VBO of  $+0.75\pm0.05$  eV was calculated for the ALD Al<sub>2</sub>O<sub>3</sub> on the bulk AlN sample, whereas a VBO of  $-0.75\pm0.05$  eV was calculated for the ALE AlN on the sapphire sample. These results are summarized in Tables I and II as well.

The REELS data for the amorphous ALD Al $_2$ O $_3$  film are shown in Fig. 5, indicating a 6.9 eV bandgap. The bandgap of crystalline sapphire substrates has been reported in several works. <sup>6.7</sup> In this case, we assumed a bandgap of 8.7 eV for the a-Al $_2$ O $_3$  substrate. <sup>7.33</sup> The difference in bandgap between amorphous and crystalline Al $_2$ O $_3$  is due to their significant difference in density (4.0 g/cm $^3$  for  $\alpha$ -Al $_2$ O $_3$  and 3.1–3.3 g/cm $^3$  for amorphous phase Al $_2$ O $_3$ ), which in turn results in a significant coordination difference. <sup>34–39</sup> The bandgap of ALE AlN has been reported previously to be 6.2 eV. The bandgap of bulk AlN was assumed to be 6.2 eV as well, as both materials are of crystalline nature. <sup>9</sup>

The resulting energy band diagrams are shown in Fig. 6, which illustrates how conduction band offsets are obtained from the valence



**FIG. 4.** High resolution XPS spectra for (a) 1.5 nm ALD  $Al_2O_3$  on AlN and (b) 1.5 nm ALE AlN on  $Al_2O_3$  (sapphire) core delta regions. The intensity is in arbitrary units (a.u.).

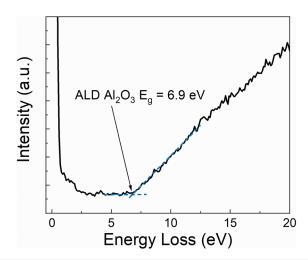


FIG. 5. Bandgap of the  ${\rm Al_2O_3}$  used in this study determined by RHEELS data. The intensity is in arbitrary units (a.u.).

band offset and the energy gap of the two constituent materials. Thus, the conduction band offset was calculated to be 1.45 eV based on the valence band offset of +0.75 eV for ALD Al<sub>2</sub>O<sub>3</sub> on bulk AlN [Fig. 6(b)] and the energy gaps of 6.9 eV and 6.2 eV for ALD Al<sub>2</sub>O<sub>3</sub> and bulk AlN, respectively, resulting in a type-II band diagram for the case of ALD amorphous Al<sub>2</sub>O<sub>3</sub> deposited on a crystalline AlN substrate. For the ALE AlN on sapphire, type-II behavior was obtained as well with a conduction band offset of 3.25 eV, owing to the larger bandgap for the sapphire substrate [Fig. 6(a)]. In both cases, the experimentally obtained conduction band offsets were significantly higher than the values reported by Robertson and Falabretti.<sup>39</sup>

Obtaining the band offsets for this heterojunction has two important consequences. First, it shows that Al<sub>2</sub>O<sub>3</sub> is a possible candidate dielectric for AlN electronic devices such as lateral AlN transistors. 10 Second, it lays the groundwork for future experiments into the band offsets for heterojunctions of ternary oxide alloys such as (Al<sub>x</sub>Ga<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub> to both binary and ternary nitrides. Similarly, our recent study of band offsets in the InN/Ga2O3 system had complemented the band offset reports for GaN/Ga2O3 and AlN/Ga2O3, allowing us to combine the data in order to estimate band offsets for the ternary nitride alloys to Ga<sub>2</sub>O<sub>3</sub>. We note that band offset data for the In<sub>2</sub>O<sub>3</sub>/GaN and In<sub>2</sub>O<sub>3</sub>/AlN heterojunctions are still lacking in the literature, as well as bowing parameter data for the (Al<sub>x</sub>In<sub>1-x</sub>)<sub>2</sub>O<sub>3</sub> ternary oxide alloy as phase segregation likely is a major issue when determining the fundamental crystal structure of this alloy. 40,41 Future work will thus focus on completing the characterization of band offsets along the ternary spectra of Al<sub>2</sub>O<sub>3</sub>, Ga<sub>2</sub>O<sub>3</sub>, AlN, GaN, In<sub>2</sub>O<sub>3</sub>, and InN.

The band offset of the Al<sub>2</sub>O<sub>3</sub>/AlN heterojunction has remained unexplored so far, as it has not been relevant for device applications. Understanding this heterostructure will pave the way for further experiments into oxide/nitride semiconductor integration. In previous studies of another potentially relevant heterojunction, InN-Ga<sub>2</sub>O<sub>3</sub>, it was proposed that knowledge of binary heterojunction band offsets could be used to calculate and predict band offsets for ternary heterojunctions.<sup>21</sup> This was done for the ternary nitride compounds AlGaN, AlInN, and InGaN with the binary compound Ga<sub>2</sub>O<sub>3</sub>. In that work, it was noted that in order to calculate the band offsets for ternary oxides to binary nitrides, knowledge of the Al<sub>2</sub>O<sub>3</sub>/AlN system is necessary. Furthermore, it was noted that eventually oxide and nitride semiconductor integration effort will be able to produce heterojunctions between ternary oxides and nitrides directly and proposed a set of simple equations to extrapolate their band offsets using the binary

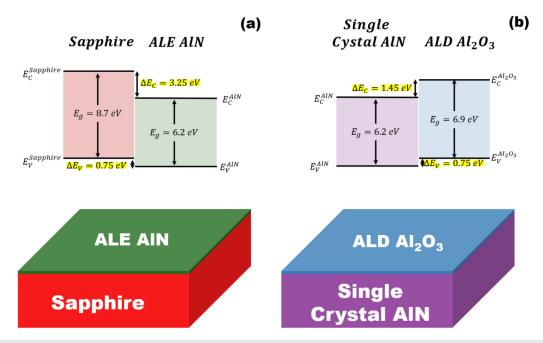


FIG. 6. Band diagrams for (a) ALD Al<sub>2</sub>O<sub>3</sub> on bulk AlN and (b) ALE AlN on a-plane sapphire (a-Al<sub>2</sub>O<sub>3</sub>)

compounds' band offset data. The measurements presented in this work present the last missing binary oxide/nitride semiconductor band offset data before any further band offset studies of ternary oxide/nitride heterojunctions can be performed.

In summary, the band alignments of ALD  $Al_2O_3$  on single crystal AlN and crystalline AlN on sapphire were measured. These are valuable in providing input to first-principles calculations of the band offsets for ternary oxide alloys such as  $(Al_xGa_{1-x})_2O_3$  and  $(In_xGa_{1-x})_2O_3$  to the binary III-nitride compounds GaN, AlN, and InN. This methodology might potentially be extended to generalize band offset calculations between heterojunctions of ternary nitrides, ternary oxides, or ternary oxide-nitride superlattices. This task is beyond the scope of this current work, primarily because additional understanding of the effects of strain, polarization, phase segregation, and interfacial defects must be developed, both experimentally and from first principles, in order to identify and minimize potential sources of error. We expect that future heterogeneous integration effort might benefit from integrating such dissimilar materials in order to meet the challenges of next-generation ultra-wide bandgap semiconductor electronic devices.

The project or effort depicted was partially sponsored by the Department of the Defense, Defense Threat Reduction Agency, No. HDTRA1-17-1-011, monitored by Jacob Calkins. The content of the information does not necessarily reflect the position or the policy of the federal government, and no official endorsement should be inferred. Research at the Naval Research Laboratory was supported by the Office of Naval Research.

### DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

## **REFERENCES**

- <sup>1</sup>J. Y. Tsao, S. Chowdhury, M. A. Hollis, D. Jena, N. M. Johnson, K. A. Jones, R. J. Kaplar, S. Rajan, C. G. Van de Walle, E. Bellotti, C. L. Chua, R. Collazo, M. E. Coltrin, J. A. Cooper, K. R. Evans, S. Graham, T. A. Grotjohn, E. R. Heller, M. Higashiwaki, M. S. Islam, P. W. Juodawlkis, M. A. Khan, A. D. Koehler, J. H. Leach, U. K. Mishra, R. J. Nemanich, R. C. N. Pilawa-Podgurski, J. B. Shealy, Z. Sitar, M. J. Tadjer, A. F. Witulski, M. Wraback, and J. A. Simmons, Adv. Electron. Mater. 4, 1600501 (2018).
- <sup>2</sup>S. J. Pearton, J. Yang, P. H. Cary, I. V. F. Ren, J. Kim, M. J. Tadjer, and M. A. Mastro, Appl. Phys. Rev. 5, 011301 (2018).
- <sup>3</sup>B. J. Baliga, Semicond. Sci. Technol. 28, 074011 (2013).
- <sup>4</sup>S. Krishnamoorthy, Z. Xia, S. Bajaj, M. Brenner, and S. Rajan, Appl. Phys. Express 10, 051102 (2017).
- <sup>5</sup>U. K. Mishra, P. Parikh, and Y.-F. Wu, "AlGaN/GaN HEMTs-an overview of device operation and applications," Proc. IEEE **90**(6), 1022–1031 (2002).
- <sup>6</sup>H. Peelaers, J. B. Varley, J. S. Speck, and C. G. Van de Walle, Appl. Phys. Lett. 112, 242101 (2018).
- <sup>7</sup>R. H. French, J. Am. Ceram. Soc. **13**, 471 (1990).
- <sup>8</sup>L. Chen and B. J. Skromme, Appl. Phys. Lett. **85**, 4334 (2004).
- <sup>9</sup>N. Nepal, S. B. Qadri, J. K. Hite, N. A. Mahadik, M. A. Mastro, and C. R. Eddy, Jr., Appl. Phys. Lett. **103**, 082110 (2013).
- 10 H. Okumura, S. Suihkonen, J. Lemettinen, A. Uedono, Y. Zhang, D. Piedra, and T. Palacios, Jpn. J. Appl. Phys., Part 1 57, 04FR11 (2018).

- <sup>11</sup>T. Kinoshita, T. Nagashima, T. Obata, S. Takashima, R. Yamamoto, R. Togashi, Y. Kumagai, R. Schlesser, R. Collazo, A. Koukitu, and Z. Sitar, Appl. Phys. Express 8, 061003 (2015).
- <sup>12</sup>Y. Taniyasu, M. Kasu1, and T. Makimoto, Nature 441, 325 (2006).
- <sup>13</sup>M. H. Breckenridge, Q. Guo, A. Klump, B. Sarkar, Y. Guan, J. Tweedie, R. Kirste, S. Mita, P. Reddy, R. Collazo, and Z. Sitar, Appl. Phys. Lett. 116, 172103 (2020).
- <sup>14</sup>R. Yamamoto, N. Takekawa, K. Goto, T. Nagashima, R. Dalmau, R. Schlesser, H. Murakami, R. Collazo, B. Monemar, Z. Sitar, and Y. Kumagai, J. Cryst. Growth 545, 125730 (2020).
- <sup>15</sup>S. Washiyama, P. Reddy, B. Sarkar, M. H. Breckenridge, Q. Guo, P. Bagheri, A. Klump, R. Kirste, J. Tweedie, S. Mita, Z. Sitar, and R. Collazo, J. Appl. Phys. 127, 105702 (2020).
- <sup>16</sup>P. Bagheri, R. Kirste, P. Reddy, S. Washiyama, S. Mita, B. Sarkar, R. Collazo, and Z. Sitar, Appl. Phys. Lett. 116, 222102 (2020).
- <sup>17</sup>M. R. Hauwiller, D. Stowe, T. B. Eldred, S. Mita, R. Collazo, Z. Sitar, and J. LeBeau, APL Mater. 8, 091110 (2020).
- <sup>18</sup>H. Ünlü and A. Asenov, J. Phys. D: Appl. Phys. 35, 591 (2002).
- <sup>19</sup>H. Sun, C. G. Torres Castanedo, K. Liu, K.-H. Li, W. Guo, R. Lin, X. Liu, J. Li, and X. Li, Appl. Phys. Lett. 111, 162105 (2017).
- <sup>20</sup>W. Wei, Z. Qin, S. Fan, Z. Li, K. Shi, Q. Zhu, and G. Z. Wei, Nanoscale Res. Lett. 7, 562 (2012).
- <sup>21</sup>C. Fares, M. J. Tadjer, J. Woodward, N. Nepal, M. A. Mastro, C. R. Eddy, Jr., F. Ren, and S. J. Pearton, ECS J. Solid State Sci. Technol. 8, Q3154 (2019).
- <sup>22</sup>T. L. Duan, J. S. Pan, and D. S. Ang, Appl. Phys. Lett. 102, 201604 (2013).
- <sup>23</sup>A. Eisenhardt, G. Eichapfel, M. Himmerlich, A. Knübel, T. Passow, C. Wang, F. Benkhelifa, R. Aidam, and S. Krischok, Phys. Status Solidi C 9(3–4), 685 (2012).
- <sup>24</sup>H. P. Song, A. L. Yang, H. Y. Weil, Y. Guol, B. Zhangl, G. L. Zheng, S. Y. Yang, X. L. Liu, Q. S. Zhu, Z. G. Wang, T. Y. Yang, and H. H. Wang, Appl. Phys. Lett. **94**, 222114 (2009).
- 25C. Fares, F. Ren, E. Lambers, D. C. Hays, B. P. Gila, and S. J. Pearton, J. Electr. Mater. 48, 1568 (2019).
- <sup>26</sup>C. Fares, M. Xian, D. J. Smith, M. R. McCartney, M. Kneiß, H. von Wenckstern, M. Grundmann, M. Tadjer, F. Ren, and S. J. Pearton, J. Appl. Phys. 127, 105701 (2020).
- <sup>27</sup>A. D. Koehler, N. Nepal, T. J. Anderson, M. J. Tadjer, K. D. Hobart, C. R. Eddy, and F. J. Kub, IEEE Electron Dev. Lett. 34, 1115 (2013).
- <sup>28</sup>P. Carey, F. Ren, D. C. Hays, B. P. Gila, S. J. Pearton, S. Jang, and A. Kuramata, Jpn. J. Appl. Phys., Part 1 56, 071101 (2017).
- <sup>29</sup>E. A. Kraut, R. W. Grant, J. R. Waldrop, and S. P. Kowalczyk, Phys. Rev. Lett. 44, 1620 (1980).
- <sup>30</sup>H. C. Shin, D. Tahir, S. Seo, Y. R. Denny, S. K. Oh, H. J. Kang, S. Heo, J. G. Chung, J. C. Lee, and S. Tougaard, Surf. Interface Anal. 44, 623 (2012).
- <sup>31</sup>E. O. Filatova and A. S. Konashuk, J. Phys. Chem. C **119**, 20755–20761 (2015).
- <sup>32</sup>A. Balzarotti and A. Bianconi, Phys. Status Solidi B **76**, 689 (1976).
- <sup>33</sup>V. V. Afanas'ev, M. Houssa, A. Stesmans, C. Merckling, T. Schram, and J. A. Kitt, Appl. Phys. Lett. **99**, 072103 (2011).
- 34X.-Y. Cui, S. P. Ringer, G. Wang, and Z. H. Stachurski, J. Chem. Phys. 151, 194506 (2019).
- 35V. A. Pustovarov, V. Sh. Aliev, T. V. Perevalov, V. A. Gritsenko, and A. P. Eliseev, J. Exp. Theor. Phys. 111, 989–995 (2010).
- 36T. V. Perevalov, V. A. Gritsenko, and V. V. Kaichev, Eur. Phys. J. Appl. Phys. 52, 30501 (2010).
- <sup>37</sup>I. Costina and R. Franchy, Appl. Phys. Lett. **78**, 4139 (2001).
- <sup>38</sup>D. Liu, S. J. Clark, and J. Robertson, Appl. Phys. Lett. **96**, 032905 (2010).
- <sup>39</sup>J. Robertson and B. Falabretti, J. Appl. Phys. **100**, 014111 (2006).
- <sup>40</sup>A. F. M. Anhar Uddin Bhuiyan, Z. Feng, J. M. Johnson, H.-L. Huang, J. Sarker, M. Zhu, M. R. Karim, B. Mazumder, J. Hwang, and H. Zhao, APL Mater. 8, 031104 (2020).
- <sup>41</sup>E. A. Anber, D. Foley, A. C. Lang, J. Nathaniel, J. L. Hart, M. J. Tadjer, K. D. Hobart, S. J. Pearton, and M. Taheri, Appl. Phys. Lett. 117, 152101 (2020).